

## Application of solving a linear system: 1-D Poisson Equation

$V$  is electrostatic potential associated with  
or gravitational potential with

Discretize second derivative and get

$$h = \quad , \quad f_n = \quad , \quad n =$$
$$x_0 <$$

and we assume that we know

$$V_0 =$$

$$V_{N+1} =$$

We get

$$n = 1$$

$$n = 2$$

$$n = 3$$

:

$$n = N$$

which is a matrix equation

$$\begin{pmatrix} & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \end{pmatrix} \begin{pmatrix} \\ \\ \\ \\ \end{pmatrix} = \begin{pmatrix} \\ \\ \\ \\ \end{pmatrix}$$

or  $A \vec{x} = \vec{b}$

with  $A$  a

Best to find routine specialized for tridiagonal matrices

→ avoid

→

show A4, Q1

show code for Poisson problem

Note: In 2D we have

Can map matrix of unknowns ( ) to a

$(m, n) =$

$i =$

→  $i =$

## Eigenvalue Problems

Many physics problems, especially from QM, can be expressed as

①

$$a_{11}x_1 + a_{12}x_2 + \dots + a_{1N}x_N = \lambda x_1$$

like system of linear eq<sup>s</sup>, except  
and solutions exist only for  
- solutions are

① can be rewritten as

For non-trivial solution ( ), this implies

, for which an  $N \times N$   
matrix leads to an

which has  $N$  solutions for  $\lambda$ , i.e.,

Finding all roots for an  $N$ -degree polynomial

Rather than solving the difficult polynomial problem,  
e.g. see Jacobi method for symmetric matrices  
(which employs "Givens rotations")  
or a general method based on QR decomposition.

Let  $A_0 =$

$$A_k = Q -$$

R - upper triangular

set  $A_{k+1} =$

$A_k$  converges to an  
- eigenvalues of an upper triangular matrix  
are

LAPACK: Different routines for different types  
of matrices.

Application: Time-Independent Schrödinger  $E\psi = \hat{H}\psi$  in 1-D

(i)

multiply by  $\psi^*$ , define  $V =$

$E =$

discretize:  $\frac{d^2 \psi}{dx^2} \rightarrow$

(i) becomes

If  $\psi(x) =$  , get

$$\begin{pmatrix} \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \end{pmatrix} = E \begin{pmatrix} \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \end{pmatrix}$$

$E =$

Each eigenvalue will have an associated eigenvector  $\vec{\psi}$ , the components of which are  $\psi(x_i)$  - the eigenfunction sampled at discrete points.

Notes on A4 Q2

Show code for A4 Q2

For harmonic potential

1)  $V(x) = \frac{1}{2} m \omega^2 x^2$  , energy eigenvalues are

2)  $E_n = \hbar \omega (n - \frac{1}{2})$  ,  $n = 1, 2, \dots$

Here

$$\hbar = \quad , \quad m =$$

so  $V(x) =$

$$E_n =$$

Now,

$$V(x) =$$

$\omega$ :

$E$ :

For 3D problems, having  $N^3$  gridpoints makes the problem difficult.

Alternative: write wavefunctions as a linear combination of

$$\psi(\vec{r}) =$$

The  $\phi_\beta(\vec{r})$ 's are known functions that solve part of the problem. E.g.

Goal is to find coefficients that solve the S.E.

$$-\nabla^2 \psi + V(\vec{r})\psi = E\psi$$

multiply by  $\phi_\alpha^*$  and integrate

$$\begin{aligned} \sum_{\beta} \int d^3\vec{r} \phi_\alpha^*(\vec{r}) (-\nabla^2 + V(\vec{r})) \phi_\beta(\vec{r}) a_\beta \\ = E \sum_{\beta} \int d^3\vec{r} \phi_\alpha^*(\vec{r}) \phi_\beta(\vec{r}) a_\beta \end{aligned}$$

Define elements of two matrices

$$H_{\alpha\beta} = \int d^3\vec{r}$$

$$S_{\alpha\beta} = \int d^3\vec{r} \quad (\text{overlap matrix})$$

Schrödinger Equation becomes

$$\sum_{\beta} = E \sum_{\beta}$$

or  $\star =$  - generalized eigenvalue problem

If  $\phi$ 's are orthonormal, then and we get

If not, need to transform  $\star$ , as through a Cholesky factorization, an LU decomposition of a positive-definite Hermitian matrix  $S$ :

$$S = \underset{\substack{\nearrow \\ \text{lower triangular}}}{L} \underset{\substack{\uparrow \\ \text{upper triangular}}}{L^+} \quad (L_{ij}^+ = L_{ji}^*)$$

$\star$  becomes

$$H I \vec{a} = E L L^+ \vec{a}$$
$$H (L^+)^{-1} L^+ \vec{a} = E L (L^+ \vec{a})$$
$$L^{-1} H (L^+)^{-1} (L^+ \vec{a}) = E (L^+ \vec{a})$$